Introduction to Octopus: a real-space (TD)DFT code

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Octopus Course 2021, MPSD Hamburg

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 - Reservation: 2 GPU and 6 CPU nodes.
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Please, only submit small jobs!

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Schedule

- Day 1 (6 Sep):
 - Short presentation of the theory
 - Key features of Octopus and numerical methods used
 - Octopus basics tutorial series
- Day 2 (7 Sep):
 - Running Octopus in HPC systems: parallelization and GPUs
- Day 3-4 (8-9 Sep):
 - Solids tutorial series
 - Optical absorption tutorial series
- Day 5 (10 Sep):
 - Model systems
- Days 6-7 (13-14 Sep):
 - Maxwell tutorials
 - Free project (students choose one or more tutorials that haven't been covered yet)

Theory

A bit of underlying theory: Density Functional Theory

- Hohenberg-Kohn theorem
- Kohn-Sham system
- time-dependent DFT

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Hohenberg-Kohn theorem

One-to-one mapping:

$$v(\vec{r}) \stackrel{W}{\longleftrightarrow} n(\vec{r})$$

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- Universal functional:

$$E_{v_0}[n] = \underbrace{\langle \Psi[n] | \hat{T} + \hat{W} | \Psi[n] \rangle}_{=F[n]} + \int d^3r \, n(\vec{r}) \, v_0(\vec{r})$$

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Variational principle:

$$E_{v_0}[n_0] = E_0$$
 for ground state n_0 $E_{v_0}[n] > E_0$ for $n \neq n_0$

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Kohn-Sham scheme

• Use one-to-one mapping: define effective non-interacting system

$$v(\vec{r}) \underbrace{\longleftrightarrow}_{W} n(\vec{r}) \underbrace{\longleftrightarrow}_{W \equiv 0} v_{s}(\vec{r})$$

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Rewrite universal functional:

$$F[n] = T_s[n] + \frac{1}{2} \int d^3r \int d^3r' \, n(\vec{r}) w(\vec{r}, \vec{r}') n(\vec{r}') + E_{xc}[n] \qquad (1)$$

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Apply variational principle (Euler-Lagrange equations):
 Kohn-Sham equations

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$$\left[-\frac{1}{2} \nabla^2 + v_{\text{ext}}(\mathbf{r}) + v_{\text{H}}[n](\mathbf{r}) + v_{\text{xc}}[n](\mathbf{r}) \right] \psi_i(\mathbf{r}) = \epsilon_i \psi_i(\mathbf{r})$$
$$n(\mathbf{r}) = \sum_{i=1}^{N} |\psi_i(\mathbf{r})|^2$$

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• One-to-one mapping (for given initial state):

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- Formulation of a variational principle based on the action is more complicated (causality paradox)
- A time-dependent Kohn-Sham system can be defined in analogy to ground-state DFT

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Time-dependent Kohn-Sham equation

$$i\frac{\partial}{\partial t}\psi_i(\boldsymbol{r},t) = \left(-\frac{1}{2}\nabla^2 + v_{\text{ext}}(\boldsymbol{r},t) + v_{\text{H}}[n](\boldsymbol{r},t) + v_{\text{xc}}[n](\boldsymbol{r},t)\right)\psi_i(\boldsymbol{r})$$
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- Various numerical schemes for doing the time-propagation

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(Time-dependent) Density Functional Theory

Summary

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- The exact exchange-correlation functionals are unknown
- The xc functionals have to be approximated

The Octopus code

Purpose: simulate the dynamics of electrons and nuclei under the influence of external time-dependent fields in the framework of Time-Dependent Density Functional Theory (TDDFT)



- DFT with many functionals (from Libxc), Hartree-Fock, Hartree
- Fortran 2008, C, C++, OpenCL/CUDA and some Python and Perl.
- extensive use of mathematical libraries: BLAS/LAPACK, FFTW, GSL, etc.
- Interfaces to external libraries: libxc, libvdwxc, wannier90, berkeleygw, etc.
- Free open-source software (GNU Public License).
- Current version is 11.1.
- Framework to implement and test new ideas

Ground-state DFT calculations



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- Excited states calculations (real-time propagation, linear response)



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 - reduced density matrix functional theory (RDMFT)



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 - van der Waals interactions



Big changes going on:

Multi-system mode



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 - classical particles (nuclei)



Introduction to Octopus

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Introduction to Octopus

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 - classical particles (nuclei)
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 - Maxwell fields



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 - \bullet currently free propagation of \vec{E} and \vec{B} fields, and propagation in linear media.



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Introduction to Octobus

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 - currently free propagation of \vec{E} and \vec{B} fields, and propagation in linear media.
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- Move to object oriented design



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More on that in advanced course.



The Octopus code

- https://octopus-code.org
- https://gitlab.com/octopus-code



M.A.L. Marques, A. Castro, G. F. Bertsch, and A. Rubio, "octopus: a first-principles tool for excited electron-ion dynamics", Comput. Phys. Commun. 151, 60-78 (2003).

A. Castro, H. Appel, M.J.T. Oliveira, C.A. Rozzi, X. Andrade, F. Lorenzen, M.A.L. Marques, E.K.U. Gross, and A. Rubio, "octopus: a tool for the application of time-dependent density functional theory", *Phys. Stat. Sol. B* 243, 2465-2488 (2006).

X. Andrade, J. Alberdi-Rodriguez, D.A. Strubbe, M.J.T. Oliveira, F. Nogueira, A. Castro, J. Muguerza, A. Arruabarrena, S.G. Louie, A. Aspuru-Guzik, A. Rubio, and M.A.L. Marques, "Time-dependent density-functional theory in massively parallel computer architectures: the octopus project", J. Phys.: Cond. Matt. 24, 233202 (2012).

X. Andrade, D.A. Strubbe, U. De Giovannini, A.H. Larsen, M.J.T. Oliveira, J. Alberdi-Rodriguez, A. Varas, I. Theophilou, N. Helbig, M.J. Verstraete, L. Stella, F. Nogueira, A. Aspuru-Guzik, A. Castro, M.A.L. Marques, and A. Rubio, "Real-space grids and the Octopus code as tools for the development of new simulation approaches for electronic systems", *Phys. Chem. Chem. Phys.* 17, 31371 (2015).

N. Tancogne-Dejean, M. J. T. Oliveira, X. Andrade, H. Appel, C. H. Borca, G. Le Breton, F. Buchholz, A. Castro, S. Corni, A. A. Correa, U. De Giovannini, A. Delgado, F. G. Eich, J. Flick, G. Gil, A. Gomez, N. Helbig, H. Hübener, R. Jestädt, J. Jornet-Somoza, A. H. Larsen, I. V. Lebedeva, M. Lüders, M. A. L. Marques, S. T. Ohlmann, S. Pipolo, M. Rampp, C. A. Rozzi, D. A. Strubbe, S. A. Sato, C. Schäfer, I. Theophilou, A. Welden, and A. Rubio, "Octopus, a computational framework for exploring light-driven phenomena and quantum dynamics in extended and finite systems", *The Journal of Chemical Physics* 152, 124119 (2020)

• Real-space grid representation



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- Finite differences for the calculation of derivatives



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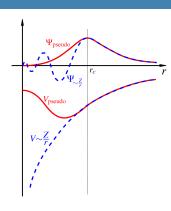
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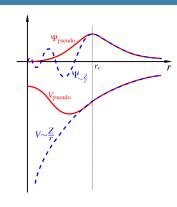
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- Quantum optimal control theory
- Many other features



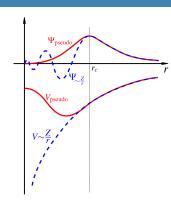
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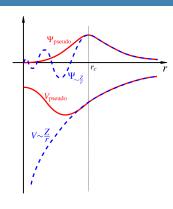
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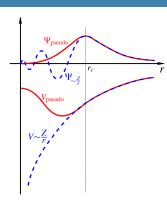
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Norm-conserving pseudo-potentials in Kleinman-Bylander form

$$V = V_{\text{loc}} + \sum_{lm} |lm\rangle (V_l - V_{\text{loc}}) \langle lm|$$

Can be specified in the Species block

• all electron potentials (delta or Gaussian nucleus)

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- user defined potentials

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- all electron potentials (delta or Gaussian nucleus)
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- and others...

Introduction to Octopus

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Introduction to Octobus

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 - (Non-uniform grids also possible)

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- Discretization: functions are represented by values on a set of points
- Point distribution:
 - Uniformly spaced grid
 - Distance between points is constant: Spacing
 - (Non-uniform grids also possible)
- Finite region of the space: Box

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 - (Decreasing spacing helps both)

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• Derivative at a point: weighted sum over neighboring points

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- General form for the Laplacian:

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Discretization of the Hamiltonian

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The Hamiltonian becomes a finite-size matrix

The eigenvalue problem

• Find the eigenvectors and eigenvalues of a matrix

The eigenvalue problem

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- Very large matrix with lots of zero components (*Sparse*)

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The eigenvalue problem

- Find the eigenvectors and eigenvalues of a matrix
- Very large matrix with lots of zero components (Sparse)
- Use iterative solvers where only the application of the matrix is required (various options available in the code)

Introduction to Octobus

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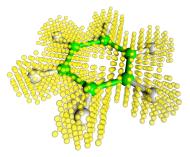
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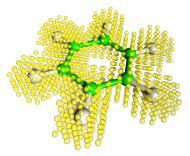
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 - etc

- Optimize the shape of the box to minimize the number of points needed
- Available box shapes:



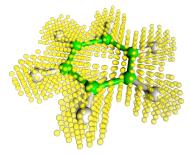
Benzene molecule in minimal box

- Optimize the shape of the box to minimize the number of points needed
- Available box shapes:
 - Sphere



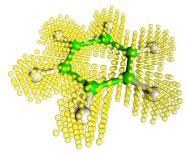
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 - Cylinder



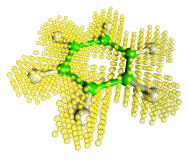
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 - Parallelepiped



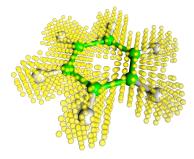
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 - Minimum box: union of spheres around each atom



Benzene molecule in minimal box

- Optimize the shape of the box to minimize the number of points needed
- Available box shapes:
 - Sphere
 - Cylinder
 - Parallelepiped
 - Minimum box: union of spheres around each atom
 - Arbitrary (e.g. 2D image!)



Benzene molecule in minimal box

Propagation of the wavefunctions in time:

$$\varphi_i(\boldsymbol{r}, t + \Delta t) = \hat{T} \exp \left\{ -\mathrm{i} \int_t^{t + \Delta t} \mathrm{d}t \, \hat{H} \varphi_i(\boldsymbol{r}, t) \right\}$$

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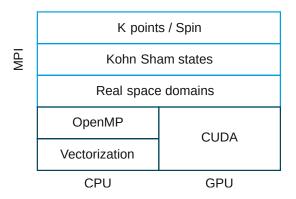
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- See e.g. tutorials on optical absorption

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Parallelisation strategy



More information in Sebastian's talk.

The tutorials

You can find the tutorials under this link: https://octopus-code.org/wiki/Tutorials

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Have fun!



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